

PH214C – Pset 6

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1 A Bohr result

Suppose that an integer number of de Broglie wavelengths fit around a circular Bohr orbit ($n \in \mathbb{Z}_+$). Find an expression for the quantized angular momentum for general n and sketch the case for $n = 4$.

$$l = mvr \quad (\text{angular momentum})$$

$$\lambda = \frac{h}{mv} \quad (\text{de Broglie wavelength})$$

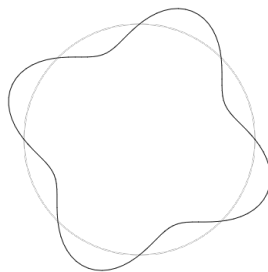
$$2\pi r = n\lambda \quad (\text{quantized orbit})$$

Thus

$$l = \frac{h}{\lambda} r = \frac{hn}{2\pi r} r = n\hbar$$

In the case where $n = 4$,

$$l_4 = 4\hbar = 4.05 \times 10^{-34} \text{ kg m}^2 \text{ s}^{-1}$$



2 Wave packet

Show that our definition of the wave function of a wave packet

$$\psi(x, t) = \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} dp \phi(p) e^{i(px - Et)/\hbar}$$

solves the free-particle Schrödinger equation

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2}$$

but not the generic wave equation:

$$\frac{\partial^2 \psi}{\partial t^2} = a^2 \frac{\partial^2 \psi}{\partial x^2}$$

2.1 Free-particle Schrödinger equation

$$i\hbar \frac{\partial \psi}{\partial t} = i\hbar \left(-\frac{iE}{\hbar} \psi \right) = E\psi \quad (\text{LHS})$$

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} = -\frac{\hbar^2}{2m} \left(\left(\frac{ip}{\hbar} \right)^2 \psi \right) = \frac{p^2}{2m} \psi = E\psi \quad (\text{RHS})$$

2.2 Generic wave equation

$$\frac{\partial^2 \psi}{\partial t^2} = \left(\frac{-iE}{\hbar} \right)^2 \psi = -\frac{E^2}{\hbar^2} \psi \quad (\text{LHS})$$

$$a^2 \frac{\partial^2 \psi}{\partial x^2} = a^2 \left(\frac{ip}{\hbar} \right)^2 \psi = -\frac{a^2 p^2}{\hbar^2} \psi \quad (\text{RHS})$$

For ψ to satisfy the wave equation, both sides must be equal:

$$-\frac{E^2}{\hbar^2} \psi = -\frac{a^2 p^2}{\hbar^2} \psi$$

$$E^2 \propto p^2$$

$$\text{but } E = \frac{p^2}{2m} \Rightarrow \text{contradiction}$$

3 Probability

We found the peak probability and expectation value $\langle x \rangle$ of the wave function

$$\psi(x) = 2a^{3/2} x e^{-ax}, \text{ for } x > 0, a = \text{constant}$$

Verify these results and then find the probability that the particle will be found between $x = 0$ and $x = 1/a$.

3.1 Peak probability

This occurs when $P = |\psi|^2$ is maximized. (Assume $a > 0$.)

$$\begin{aligned} P(x) &= |\psi(x)|^2 = 4a^3 x^2 e^{-2ax} \\ \frac{dP}{dx} &= 8a^3 x(1 - ax)e^{-2ax} = 0 \\ \frac{d^2P}{dx^2} &= 8a^3(1 - 4ax + 2a^2 x^2)e^{-2ax} \end{aligned}$$

This equation is satisfied when $x = 0$ or $x = 1/a$. Using the second derivative:

$$\begin{aligned} \left. \frac{d^2P}{dx} \right|_{x=0} &= 8a^3 > 0 \\ \left. \frac{d^2P}{dx} \right|_{x=1/a} &= -8a^3 e^{-2} < 0 \end{aligned}$$

Thus the only local maximum occurs at $x = 1/a$. (Also, since the probability dies out as $x \rightarrow \infty$, this must be the global maximum in the interval $[0, \infty)$.)

3.2 Expectation value

Since $\psi = 0$ for negative x , we can begin the integration at $x = 0$. Since ψ and x are real, $P = |\psi|^2 = \psi^* \psi = \psi^2$, and $\psi^* x \psi = x \psi^2 = xP$.

$$\begin{aligned} \langle x \rangle &= \int_{-\infty}^{\infty} dx \psi^* x \psi \\ &= \int_0^{\infty} dx x (4a^3 x^2 e^{-2ax}) \\ &= \frac{1}{4a} \int_0^{\infty} dy y^3 e^{-y} \quad (y = 2ax) \\ &= \frac{1}{4a} \left[e^{-y}(-y^3 - 3y^2 - 6y - 6) \right]_{x=0}^{\infty} \quad (\text{a lot of integration by parts}) \\ &= \frac{3}{2a} \end{aligned}$$

3.3 Probability of position between 0 and $1/a$

This is just an integral of probability over some interval of x .

$$\begin{aligned}
 P\left(0 \leq x \leq \frac{1}{a}\right) &= \int_0^{1/a} dx P(x) \\
 &= \int_0^{1/a} dx (4a^3 x^2 e^{-2ax}) \\
 &= \frac{1}{2} \int_0^2 dy y^2 e^{-y} \quad (y = 2ax) \\
 &= \frac{1}{2} \left[e^{-y}(-y^2 - 2y - 2) \right]_0^2 \quad (\text{integration by parts}) \\
 &= \frac{1}{2} (e^{-2}(-4 - 4 - 2) - e^0(-2)) \\
 &= 1 - 5e^{-2} \approx 0.32
 \end{aligned}$$

4 Momentum keeps it real

Show that $\langle p \rangle - \langle p \rangle^* = 0$ to show that the operator for p is Hermitian.

$$\begin{aligned}
 \hat{p} &:= \frac{\hbar}{i} \frac{\partial}{\partial x} \\
 \langle p \rangle - \langle p \rangle^* &= \left[\int_{-\infty}^{\infty} dx \psi^* \left(\frac{\hbar}{i} \frac{\partial}{\partial x} \right) \psi \right] - \left[\int_{-\infty}^{\infty} dx \psi^* \left(\frac{\hbar}{i} \frac{\partial}{\partial x} \right) \psi \right]^* \\
 &= \frac{\hbar}{i} \left[\int_{-\infty}^{\infty} dx \psi^* \frac{\partial \psi}{\partial x} \right] - \left(\frac{\hbar}{-i} \right) \left[\int_{-\infty}^{\infty} dx \psi \frac{\partial \psi^*}{\partial x} \right] \\
 &= \frac{\hbar}{i} \int_{-\infty}^{\infty} dx \psi^* \frac{\partial \psi}{\partial x} + \psi \frac{\partial \psi^*}{\partial x} \\
 &= \frac{\hbar}{i} \int_{-\infty}^{\infty} dx \frac{\partial}{\partial x} (\psi^* \psi) = \frac{\hbar}{i} \int_{-\infty}^{\infty} dx \frac{\partial |\psi|^2}{\partial x} \\
 &= \frac{\hbar}{i} \lim_{b \rightarrow \infty} |\psi|^2 \Big|_{x=-b}^b = \frac{\hbar}{i} [0 - 0] = 0
 \end{aligned}$$

Notes about the math:

- Any valid PDF (e.g., $|\psi|^2$) must approach the value 0 in the limit as the variable goes to $\pm\infty$.
- In general,

$$\left[\int f(x) dx \right]^* = \int [f(x)]^* [dx]^*$$

Since x and dx are real, then the conjugate of the integral is the integral of the conjugate of the integrand.