## PHYSICAL CHEMISTRY

Chapter 5: The Second Law of Thermodynamics

#### 5.1. Spontaneous Processes

- The reverse of a spontaneous process won't happen under the same conditions
- Viewing events with the random movement of many molecules makes certain events very unlikely to occur (i.e., these events would require the synchronization of a huge number of molecules)
- Cannot base spontaneity on a system losing energy (although this is common for spontaneous events), but rather entropy

Significant examples:

• Ball bouncing and losing thermal energy with each bounce; reverse process (gaining thermal energy from ground and rising up) cannot be spontaneous (intuitively)

### 5.2. Entropy

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- <u>Statistical definition of entropy</u>: entropy is proportional to the number of molecules N, and probability is related to an exponential function of N, so let entropy  $S \propto \ln W$ , S is entropy, W is probability
  - Let the proportionality constant be  $k_B$  (Boltzmann's constant), so  $S = k_B \ln W$
  - Because S is a state function, the path doesn't matter
  - For calculating changes in entropy:  $\Delta S = S_2 S_1 = k_B \ln \frac{W_2}{W_1}$
  - For a gas,  $W = (CV_1)^N$  (approximately volume raised to the *N*th number of molecules), so  $\Delta S = k_B \ln \left(\frac{V_2}{V_1}\right)^N = nR \ln \frac{V_2}{V_1}$ 
    - Only for isothermal expansion, because entropy is affected by temperature

• Thermodynamic definition of entropy: 
$$\Delta S = \frac{q_{rev}}{T}$$
 (or  $dS = \frac{dq_{rev}}{T}$ )

- Derivation: For an isothermal, reversible expansion:  $q = -w; q_{rev} = nRT \ln \frac{V_2}{V_1} \Rightarrow \frac{q_r ev}{T} = nR \ln \frac{V_2}{V_1} = \Delta S$
- Note that this equation works for all substances (not only gases)
- Note that this equation requires the reversible q (even though  $\Delta S_{rev} = \Delta S_{irrev}$  because state function, this calculation requires  $q_{rev}$ )

Significant examples:

• Likelihood of finding N molecules in half a container vs. the whole container

### 5.3. The Carnot Heat Engine

- A <u>heat engine</u> converts heat to mechanical work (e.g., steam engines and steam-powered electricity generators)
- <u>Carnot heat engine</u> is an idealized heat engine of one mole of gas with four steps:
  - Four steps:
    - Isothermal, reversible expansion from  $V_1$  to  $V_2$  (at  $T_2$ )
    - Adiabatic, reversible expansion from  $V_2$  to  $V_3$  ( $T_2 \rightarrow T_1$ )
    - Isothermal, reversible compression from  $V_3$  to  $V_4$  (at  $T_1$ )
    - Adiabatic, reversible compression from  $V_4$  to  $V_1$  ( $T_1 \rightarrow T_2$ )
  - **Properties:** 
    - $\Delta U(\text{cycle}) = 0$

- $q(\text{cycle}) = RT_2 \ln \frac{V_2}{V_1} + RT_1 \ln \frac{V_4}{V_2}$ 

  - $q_2 = RT_2 \ln \frac{V_2}{V_1}$  (heat absorbed from heat reservoir, positive value)  $q_1 = RT_1 \ln \frac{V_4}{V_3} = -RT_1 \ln \frac{V_2}{V_1}$  (heat discharged to cold reservoir, negative value)

• 
$$w(\text{cycle}) = -RT_2 \ln \frac{V_2}{V_1} + \bar{C}_V(T_1 - T_2) - RT_1 \ln \frac{V_4}{V_3} + \bar{C}_V(T_2 - T_1) = -RT_2 \ln \frac{V_2}{V_1} - RT_1 \ln \frac{V_4}{V_3}$$

- From volume equation below,  $w(\text{cycle}) = -R(T_2 T_1) \ln \frac{V_2}{V_1}$
- $\frac{V_2}{V} = \frac{V_3}{V}$

• Derivation: 
$$P_1V_1 = P_2V_2$$
;  $P_3V_3 = P_4V_4$ ;  $P_2V_2Y^{\gamma} = P_3V_3^{\gamma}$ ;  $P_1V_1^{\gamma} = P_4V_4^{\gamma}$ ;  
 $\frac{P_2V_2^{\gamma}}{P_1V_1^{\gamma}} = \frac{P_3V_3^{\gamma}}{P_4V_4^{\gamma}} \Rightarrow \frac{P_2V_2}{P_1V_1} \times \left(\frac{V_2}{V_1}\right)^{\gamma-1} = \frac{P_3V_3}{P_4V_4} \times \left(\frac{V_3}{V_4}\right)^{\gamma-1} \Rightarrow \left(\frac{V_2}{V_1}\right)^{\gamma-1} = \left(\frac{V_3}{V_4}\right)^{\gamma-1}$ 

Thermodynamic <u>efficiency</u> is the ratio of work done by heat engine to heat absorbed by engine ( efficiency =  $\frac{\text{net work done by heat engine}}{\text{heat absorbed by engine}} = \frac{|w|}{a}$ )

• For Carnot engine, efficiency = 
$$\frac{R(T_2 - T_1) \ln \frac{V_2}{V_1}}{RT_2 \ln \frac{V_2}{V_1}} = \frac{T_2 - T_1}{T_2} = 1 - \frac{T_1}{T_2}$$

• Cannot be totally efficient; efficiency maximized by maximizing difference between *T*<sub>1</sub> and  $T_2$ 

• For any cyclic process, 
$$\sum_{i} \frac{q_i}{T_i} = \sum_{i} \Delta S_i = 0$$

- Refrigerators, air conditioners, and heat pumps reverse the flow of spontaneous heat flow by applying work:  $-q_2 = q_1 + w$ 
  - Coefficient of performance (COP) is measure of refrigerator or air conditioner's performance:  $\text{COP} = \frac{q_1}{w} = \frac{T_1}{T_2 - T_1}$ ; gives the maximum COP value because it uses reversible heat value
  - COP for heat pumps is  $\frac{T_2}{T_2 T_1}$ ; more efficient than an electric heater

### 5.4. The Second Law of Thermodynamics

Total change in entropy for a reversible process,  $\Delta S_{surr} = 0$ ,  $q_{surr} = -q_{sys}$ • For an irreversible process,  $\Delta S_{univ} \geq 0$ 

# 5.5. Entropy Changes

- Entropy change for reversible, isothermal expansion of an ideal gas is  $nR \ln \frac{V_2}{V_2}$ ٠
- $\Delta_{mix}S = n_A R \ln \frac{V_A + V_B}{V_A} + n_B R \ln \frac{V_A + V_B}{V_B}$  for mixing of two gases (can be thought of as

two gas expansions)  $\circ \Delta_{mix} S_{A} = -\frac{R}{rr} (n_A \ln x_A + n_B \ln x_B)$  (greater than 0, so spontaneous)

• 
$$\Delta_{fus}S = \frac{\Delta_{fus}H}{T_f}$$
 (because enthalpy of fusion is equal to reversible heat)

• Trouton's rule is an empirical observation that many liquids have a similar entropy of vaporization (because of similar structures of most liquids and gases)

•  $\Delta S = C_p \ln \frac{T_2}{T_1}$  when heating